

# On “quasi-Richards” equation and finite volume approximation of two-phase flow with unlimited air mobility

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Passage to the limit (singular) and “quasi-Richards”

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  - Passage to the limit (singular) and “quasi-Richards”
  - Renormalization and incomplete contraction inequality
- 3 Finite volumes for two-phase flow with unlimited mobility
  - The idea of the scheme
  - A priori estimates, existence, convergence at fixed  $\mu$
  - Asymptotics of the scheme as  $\mu \rightarrow 0$
  - Numerical illustrations

## Models

### Assumptions about groundwater flow

Water and air      incompressible phases

Porous medium    homogeneous and isotropic

Gravity            neglected

Source term      of a special form

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Water and air	incompressible phases	
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### Richards model

$$\begin{cases} u_t - \operatorname{div}(k_w(u)\nabla p) = s_w, \\ u = p_c^{-1}(p_{atm} - p), \end{cases}$$

### Two-phase model

$$\begin{cases} u_t - \operatorname{div}(k_w(u)\nabla p) & = s_w \\ (1 - u)_t - \operatorname{div}(\mu k_a(u)\nabla(p + p_c(u))) & = s_a \end{cases}$$

where  $\mu :=$  Ratio between the phase mobilities (we want  $\mu \rightarrow \infty$ )

## Assumptions

- 1  $\Omega$  is a polygonal subset of  $\mathbb{R}^d$ ,  $d = 2$  or  $3$ ;  $T > 0$  is given,
- 2  $u_m \in (0, 1)$ ; initial saturation  $u_0$ , source saturation  $c$  with  $u_m \leq u_0(x) \leq 1$  a.e on  $\Omega$ ,  
and  $u_m \leq c(t, x) \leq 1$  a.e. on  $\Omega \times (0, T)$ ,
- 3 source  $\bar{s} \in L^2$ , sink  $\underline{s} \in L^2$ ,  $\bar{s}, \underline{s} \geq 0$ ,  
global conservation:  $\int_{\Omega} (\bar{s}(x, t) - \underline{s}(x, t)) dx = 0$  on  $(0, T)$ ,
- 4  $k_w \in C^0([0, 1])$ ,  $k_w$  non-decreasing with  $k_w(0) = 0$ ,  
 $k_w(1) = 1$  and  $k_w(u_m) > 0$ ,  
 $k_a \in C^0([0, 1])$ ,  $k_a$  non-increasing with  $k_a(1) = 0$ ,  
 $k_a(0) = 1$  and  $k_a(s) > 0$  for all  $s \in [0, 1)$ ,  
 $p_c \in C^0([u_m, 1]) \cap \text{Lip}_{loc}([u_m, 1])$ ,  $p_c$  strictly decreasing
- 5  $\mu \in [1, +\infty)$

## Mathematical setting of the problem

Set  $f_\mu(u) := \frac{k_w(u)}{k_w(u) + \mu k_a(u)}$ ; one has  $f_\mu \xrightarrow{\mu \rightarrow \infty} \mathbb{1}_{[u=1]}$

**Two-phase problem: find  $(u, p)$  such that:**

$$\left\{ \begin{array}{l} u_t - \operatorname{div}(k_w(u) \nabla p) = f_\mu(c) \bar{s} - f_\mu(u) \underline{s} \quad \text{on } \Omega \times (0, T), \\ \end{array} \right.$$

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To let  $\mu \rightarrow \infty$ : uniform estimates for discrete/regularized problem

## The “quasi-Richards” equation

### Theorem (Eymard, Henry, Hilhorst'09, DCDS-S'12.)

There exist solutions  $(u^\mu, p^\mu)$  for the two-phase flow problem that obey *uniform estimates*: lower bound  $u_m$  on the saturations  $u^\mu$ ,  $L^2(0, T; H^1)$  bound on the pressures  $p^\mu$  and on the **1/2-Kirchoff transform**  $\zeta(u^\mu)$ ,

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$$\begin{cases} u_t - \operatorname{div}(k_w(u)\nabla p) = s_w \\ \nabla(p + p_c(u)) = 0 \text{ a.e. on the set } [u < 1], \end{cases}$$

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Thus: solution of the quasi-Richards eqn. is a **triple**  $(u, p, \theta)$  with  $\nabla p = -\nabla p_c(u)$  on  $[u < 1]$  and with  $\theta$  defined on  $[u = 1]$ .

Regularity:

$u$  is  $[u_m, 1]$ -valued with  $\zeta(u) \in L^2(0, T; H^1)$ ,  $p \in L^2(0, T; H^1)$ .

## Formal justification of the estimates:

- Multiply first equation by  $p$ , second equation by  $(p + p_c(u))$ , sum up, use chain rule on  $p_c(u)u_t$ , use  $k_w(u) \geq k_w(u_m) > 0$   
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- Write a “global flux” formulation:

$$-\operatorname{div} \mathbf{q} = \bar{s} - \underline{s}, \quad \mathbf{q} \cdot \mathbf{n} = 0,$$

$$u_t - \operatorname{div}[f_\mu(u)\mathbf{q} - k_w(u)\nabla Q(u)] = s_w, \quad \nabla Q(u) \cdot \mathbf{n} = 0$$

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Eliminate  $q$  from the resulting system: test fct  $-p_c(u)$  in the 2nd, test function  $F_\mu(u) := \int_0^u f_\mu(s) p'_c(s) ds$  in the 1st equation  $\Rightarrow$

$$k_w(u) \nabla Q(u) \cdot \nabla p_c(u) = \frac{k_a(u) \mu k_w(u)}{k_w(u) + \mu k_a(u)} |p'_c(u)|^2 |\nabla u|^2 \text{ bounded in } L^1$$

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$\Rightarrow$  using  $k_w(u) \geq k_w(u_m) > 0$ , we get  $L^1$  bound on

$$|\nabla \zeta(u)|^2 = k_w(u) |p'_c(u)| |\nabla u|^2 \geq \frac{k_a(u) \mu k_w(u)}{k_w(u) + \mu k_a(u)} |p'_c(u)|^2 |\nabla u|^2.$$

## Justification of passage to the limit; multiplier $\theta$

- In addition to estimate of  $\nabla\zeta(u^\mu)$ , use translation estimates in time to get strong compactness of  $\zeta(u^\mu)$ ; use invertibility of  $\zeta(\cdot)$   
 $\Rightarrow$  create a strong limit  $u$  of  $u^\mu$  .

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- Pass to strong limit in nonlinearities  $k_a(u^\mu)$ ,  $k_w(u^\mu)$ , to weak limit in  $\nabla p^\mu \Rightarrow$  create a weak limit  $p$  of  $p^\mu$  .

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- On the set  $[u < 1]$  where  $k_a(u) > 0$ ,  
 pass to the limit  $\mu \rightarrow \infty$  in  $L^1$  bound  $k_a(u^\mu)\mu|\nabla(p + p_c(u^\mu))|^2$   
 $\Rightarrow$  get 2nd line (constraint) of weak quasi-Richards formulation .

## Is quasi-Richards well-posed? Is it different from Richards?

- Richards is well-posed: [Alt, Luckhaus'83](#) .  
 $L^1$  contraction inequality holds. ( $\Rightarrow$  uniqueness, stability)
- Existence of sols to quasi-Richards: [Eymard, Henry, Hilhorst](#) .  
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### Theorem ([A., Eymard, Ghilani, Marhraoui'12](#))

Assume  $u, \hat{u}$  are weak solutions of the quasi-Richards equation corresponding to data  $(u_0, \bar{s})$  and  $(\hat{u}_0, \widehat{\bar{s}})$ . Then we have the following [incomplete contraction inequality](#) : for a.e.  $t$ ,

$$\int_{\Omega} (u - \hat{u})^+(t, \cdot) \leq \int_{\Omega} (u_0 - \hat{u}_0)^+ + \int_0^t \int_{\Omega} (\bar{s} - \widehat{\bar{s}})^+ + \int_0^t \int_{[u=1=\hat{u}]} \bar{s}. \quad (1)$$

Proof: [use renormalized solutions](#) of [Plouvier-Debaigt, Gagneux](#) .

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### Theorem (A., Eymard, Ghilani, Marhraoui'12 )

*Assume there is no water injection:  $\bar{s}\mathbb{1}_{[c=1]} = 0$  a.e. on  $(0, T) \times \Omega$  (with  $c = c(t, x)$  the saturation in water of the injected fluid).*

*Then for every datum  $u_0$  there exists a unique  $u$  such that  $(u, p, \theta)$  is a solution of the quasi-Richards equation.*

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### Theorem (A., Eymard, Ghilani, Marhraoui'12 )

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*Then for every datum  $u_0$  there exists a unique  $u$  such that  $(u, p, \theta)$  is a solution of the quasi-Richards equation.*

*Moreover, in absence of water injection we have  $\theta \underline{s} = 0$  a.e. (no water production!); and the saturation  $u$  given by quasi-Richards eqn coincides with the unique solution of the Richards eqn.*

## Is quasi-Richards well-posed? Is it different from Richards?

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In general, we do not expect that quasi-Richards and Richards coincide:

- Physical reasons:  $p_{atm}$  is not the good pressure for air when air is captured by saturated water phase
- While uniqueness of  $u$  in the triple  $(u, p, \theta)$  can be hoped for, we do not expect uniqueness of  $(p, \theta)$  in the saturated set  $[u = 1]$ .

More work needed to understand quasi-Richards !

## Renormalized solutions...

- Idea: multiply quasi-Richards by a nonlinear truncation  $T_n(u)$ ,

$$T_n \equiv 1 \text{ on } [0, 1 - \frac{1}{n}], \quad T_n(1) = 0.$$

Use chain rules, obtain family of evolution equations

$$(RenEq_n) \quad b_n(u)_t - \Delta \varphi_n(u) = \bar{s} T_n(u) + |\nabla \psi_n(u)|^2$$

with *ad hoc* nonlinearities  $b_n, \varphi_n, \psi_n$ .

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$$(Cstr) \quad \lim_{n \rightarrow \infty} \int_0^t \int_{\Omega} |\nabla \psi_n(u)|^2 = \int_0^t \int_{[u=1]} (\bar{s} - \theta \underline{s}).$$

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Def. Combination of equations  $(RenEq_n)$ ,  $n \rightarrow \infty$  and of constraint  $(Cstr)$  is called **renormalized formulation** of quasi-Richards.

Prop. A weak solution is also a renormalized solution.

## Use of renormalized solutions.

- ( $RenEq_n$ ) is a “standard parabolic-elliptic problem”  $\Rightarrow$   $L^1$ -contraction ok. Given two solutions  $u, \hat{u}$ , we find

$$\begin{aligned} & \| (b_n(u) - b_n(\hat{u}))^+ \|_{L^1}(t) \leq \| (b(u_0) - b(\hat{u}_0))^+ \|_{L^1} \\ & + \int_0^t \int_{\Omega} \text{sign}^+(b_n(u) - b_n(\hat{u})) \left( \bar{s} T_n(u) - \hat{s} T_n(\hat{u}) + |\nabla \psi_n(u)|^2 - |\nabla \psi_n(\hat{u})|^2 \right). \end{aligned}$$

- Let  $n \rightarrow \infty$  : using  $b_n \rightarrow \text{Id}$ , we get

$$\begin{aligned} & \| (u - \hat{u})^+ \|_{L^1}(t) \leq \| (u_0 - \hat{u}_0)^+ \|_{L^1} \\ & + \int_0^t \int_{[u > \hat{u}]} \left( \bar{s} \mathbb{1}_{[u < 1]} - \hat{s} \mathbb{1}_{[\hat{u} < 1]} \right) + \lim_{n \rightarrow \infty} \int_0^t \int_{u > \hat{u}} \left( |\nabla \psi_n(u)|^2 - |\nabla \psi_n(\hat{u})|^2 \right). \end{aligned}$$



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- Use ( $Cstr$ ) trying to simplify the right-hand side...

...aie aie... the term  $\int_0^t \int_{[u=1=\hat{u}]} \bar{s}$  survives.

$\Rightarrow$  incomplete contraction inequality follows.

## Finite volume scheme

Write  $k_w(u) = f_\mu(u) M_\mu(u)$ ,  $M_\mu = k_w + \mu k_a$ . Set  $\delta_{K,L}^{n+1}(Z_D) = Z_L^{n+1} - Z_K^{n+1}$ .

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 The scheme is: find  $U_D = (U_K^n)_{n,K}$ ,  $P_D = (P_K^n)_{n,K}$  satisfying

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- $U_{K|L}^{n+1}$  is the upwind value : 
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- $\bar{U}_{K|L}^{n+1}$  between  $U_K^{n+1}$  and  $U_L^{n+1}$  is the auxiliary value:

$$k_a(\bar{U}_{K|L}^{n+1}) \delta_{K,L}^{n+1}(p_c(U_D)) = \delta_{K,L}^{n+1}(g(U_D)) \text{ i.e., } k_a(\bar{U}_{K|L}^{n+1}) = \frac{g(U_L^{n+1}) - g(U_K^{n+1})}{p_c(U_L^{n+1}) - p_c(U_K^{n+1})}.$$

## Properties of the scheme

- the choice  $\bar{U}_{K|L}^{n+1}$  makes appear  $\mu k_a(\bar{U}_{K|L}^{n+1}) \delta_{K,L}^{n+1} (P_{\mathcal{D}} - p_c(U_{\mathcal{D}})) \Rightarrow$  uniform in  $\mu, h$  (discrete) estimates as for [Eymard, Henry, Hilhorst](#) ... except for time translation estimate on  $U_{\mathcal{D}}$  (not uniform in  $\mu$ )

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In the gradually saturated regime ( $u \leq u_M < 1$ ) we find

$$\frac{U_K^{n+1} - U_K^n}{\delta t^n} m_K - \sum_{L \in \mathcal{N}_K} \tau_{K|L} k_w(U_{K|L}^{n+1}) \frac{k_a(\bar{U}_{K|L}^{n+1})}{k_a(U_{K|L}^{n+1})} \delta_{K,L}^{n+1} (P_D) = 0,$$

while the straightforward discretization of Richards equation yields

$k_w(U_{K|L}^{n+1}) \delta_{K,L}^{n+1} (P_D)$ . One can see that  $\frac{k_a(\bar{U}_{K|L}^{n+1})}{k_a(U_{K|L}^{n+1})} \rightarrow 1$  as  $h \rightarrow 0$ , so we

have an “almost asymptotic preserving” scheme: (limit  $\mu \rightarrow \infty$  of the two-phase scheme is a “strange” scheme for Richards eqn).

Merci pour votre attention !