
Strong boundary traces and well-posedness for scalar conservation laws with dissipative boundary conditions

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1 Introduction

The aim of this paper is to give sense to the following formal problem for a scalar conservation law with boundary condition (BC, in the sequel) :

$$(H) \begin{cases} u_t + \operatorname{div} \varphi(u) = f & \text{in } Q := (0, T) \times \Omega \\ u(0, \cdot) = u_0 & \text{on } \Omega \\ \varphi_\nu(u) := \varphi(u) \cdot \nu \in \beta(u) & \text{on } \Sigma := (0, T) \times \partial\Omega. \end{cases}$$

Here $\Omega = \mathbb{R}^+ \times \mathbb{R}^{N-1}$ ($N \geq 1$), $T > 0$, ν is the unit outward normal vector on $\partial\Omega$, $\varphi : \mathbb{R} \rightarrow \mathbb{R}^N$ is continuous, and β is a maximal monotone graph on \mathbb{R} . Assume $\varphi_\nu(0) - \beta(0) \ni 0$ and normalize φ, β by $\varphi(0) = 0$, $0 \in \beta(0)$. The classical Neumann (zero-flux) and Dirichlet homogeneous BC correspond to the graphs $\beta = \mathbb{R} \times \{0\}$ and $\beta = \{0\} \times \mathbb{R}$, respectively.

To show existence, we restrict our attention to the case of Ω with flat boundary, to $L^1 \cap L^\infty$ data u_0, f (see (12) for the precise assumption on the data), and make the following simplifying assumptions:

$$\text{there exists a constant } C \text{ such that } |\beta(z)| \geq \operatorname{sign}(z) \varphi_\nu(z) \quad \forall |z| > C; \quad (1)$$

$$\varphi \text{ is Lipschitz continuous, and } \varphi_\nu = \varphi \cdot \nu \text{ is piecewise monotone;} \quad (2)$$

$$\forall \xi \in \mathbb{R}^N \setminus \{0\}, \text{ the function } z \mapsto \xi \cdot \varphi(z) \text{ is non-constant on any interval.} \quad (3)$$

In the Conclusion, possible generalizations are indicated.

It is well known that existence for (H) generally fails if one interprets the Dirichlet BC literally (cf. [BLN]). This is also the case for general β , except for some particular situations (cf. e.g. [BFK]). If one approximates (H) by a sequence of problems (H^ϵ) with the BC understood literally (e.g. parabolic “viscous” approximations, or numerical schemes), a boundary layer can form

in the corresponding solutions u^ϵ . The convergence of u^ϵ then takes place only locally inside the domain, and the limiting function u satisfies to the scalar conservation law with a different BC, which we call “effective BC”.

It is the goal of this paper to investigate the form of the effective BC for a general BC given by a maximal monotone graph β . Note that the monotonicity is necessary, if we hope the boundary condition to be dissipative in L^1 (in particular, if we hope for an L^1 contraction principle for solutions of (H), as in the classical situations of $\Omega = \mathbb{R}^N$, or of bounded Ω with Dirichlet BC).

We suggest that the effective BC is given by a monotone graph $\tilde{\beta}$ defined by

$$\tilde{\beta} := \left\{ \begin{array}{l} \exists b \in \overline{\text{Range}(\beta)} \text{ such that } \varphi_\nu(z) = b \text{ and} \\ (z, \varphi_\nu(z)); \text{ if } z < m := \inf \beta^{-1}(b), \text{ then } \varphi_\nu(k) \geq b \ \forall k \in [z, m[\\ \text{if } z > M := \sup \beta^{-1}(b), \text{ then } \varphi_\nu(k) \leq b \ \forall k \in]M, z] \end{array} \right\}, \quad (4)$$

which can be visualized as the horizontal projection of β on the graph of φ_ν . The graph $\tilde{\beta}$ is monotone, and $\tilde{\tilde{\beta}} = \tilde{\beta}$ (thus operation $\tilde{\cdot}$ is indeed a projection).

Example 1. (i) If $\beta = \{0\} \times \mathbb{R}$, then $\tilde{\beta}$ is the Bardos-LeRoux-Nédélec graph:

$$\tilde{\beta} = \{(z, \varphi_\nu(z)) \mid \text{sign}(z)(\varphi_\nu(z) - \varphi_\nu(k)) \geq 0 \ \forall k \in [0 \wedge z, 0 \vee z]\}. \quad (5)$$

(ii) If $\beta = \mathbb{R} \times \{0\}$, then $\tilde{\beta} = \{(z, \varphi_\nu(z)) \mid \varphi_\nu(z) = 0\}$. Assumption (1) is restrictive in this case; a similar assumption is made in [BFK]. For the general case, one has to complete β to a maximal monotone graph on $[-\infty, +\infty]$ before defining $\tilde{\beta}$ by (4); see Conclusion and the forthcoming paper [AS].

2 Strong boundary traces for entropy solutions

The Bardos-LeRoux-Nédélec pointwise formulation of the Dirichlet BC, which can be expressed by means of the graph (5) was initially given for BV solutions (compare to the work of Otto in [MNR], where a weak formulation of the Dirichlet boundary condition is given, valid for any L^∞ solution). It has recently been realized that any L^∞ entropy solution actually has strong L^1_{loc} initial and boundary traces in a fairly general situation (see Panov [P05, P06] and the previous works of Chen-Rascle and Vasseur). In particular, the non-degeneracy condition (3) is not needed. The concept of strong trace (in L^1_{loc}) is stated in the following definition. Set $x = (x_1, x')$, $x_1 \in \mathbb{R}^+$, $x' \in \mathbb{R}^{N-1}$.

Definition 1. A function $\tilde{v} \in L^1_{loc}(\mathbb{R}^{N-1})$ is a strong trace of function $v \in L^1_{loc}(\mathbb{R}^+ \times \mathbb{R}^{N-1})$ on $\{x_1 = 0\}$ if for all $\xi \in C_c(\mathbb{R}^{N-1})$, $\xi \geq 0$

$$\text{ess-} \lim_{x_1 \rightarrow 0} \int_{\mathbb{R}^{N-1}} \xi(x') |v(x_1, x') - \tilde{v}(x')| dx' = 0.$$

In the sequel $\gamma : L^1_{loc}(\mathbb{R}^+ \times \mathbb{R}^{N-1}) \rightarrow L^1_{loc}(\mathbb{R}^{N-1})$ is the strong trace operator in the sense of Definition 1. Traces of $L^1_{loc}(Q)$ functions on $\Sigma = (0, T) \times \mathbb{R}^{N-1}$ and on $\{t = 0\}$ are defined similarly. Clearly, the strong trace operators are unbounded. The following proposition, inferred by the result of Panov [P05] on traces of entropy (quasi-)solutions at $\{t = 0\}$, is therefore remarkable.

Proposition 1. *Assume (2). Let u be a quasi-solution for $u_t + \operatorname{div} \varphi(u)^3$. Then there exists a strong trace $\gamma V_{\varphi_\nu}(u)$ of the function $V_{\varphi_\nu}(u)$ on Σ in the sense of Definition 1, where $V_{\varphi_\nu}(\cdot)$ is the variation function of φ_ν :*

$$V_{\varphi_\nu}(z) = \operatorname{sign}(z) \operatorname{Var}_{[0 \wedge z, 0 \vee z]} \varphi_\nu(\cdot).$$

In particular, $\operatorname{sign}^\pm(u - k)(\varphi(u) - \varphi(k)) \cdot \nu$, which can be rewritten as

$$Q^\pm(V_{\varphi_\nu}(u), V_{\varphi_\nu}(k)) := \operatorname{sign}^\pm(V_{\varphi_\nu}(u) - V_{\varphi_\nu}(k)) (\Psi_\nu(V_{\varphi_\nu}(u)) - \Psi_\nu(V_{\varphi_\nu}(k)))$$

with $\Psi_\nu := \varphi_\nu \circ V_{\varphi_\nu}^{-1}$, has the strong trace $Q^\pm(\gamma V_{\varphi_\nu}(u), V_{\varphi_\nu}(k))$ on Σ .

Notice that the above result does not depend on (1),(3). It still holds if Ω is a domain with piecewise C^1 -smooth boundary and $\varphi \in BV$ (see Panov [P06]).

3 Entropy solutions and well-posedness

Definition 2. *A function $u \in L^\infty(Q)$ is said an entropy solution for Problem (H) if $\forall k \in \mathbb{R}$, $\forall \xi \in C_c^\infty(Q)$, $\xi \geq 0$ the local entropy inequalities hold :*

$$\int_Q (u - k)^\pm \xi_t + \int_Q \operatorname{sign}^\pm(u - k)(\varphi(u) - \varphi(k)) \cdot D\xi + \int_Q f \operatorname{sign}^\pm(u - k) \xi \geq 0, \quad (6)$$

u has the strong trace u_0 on $\{t = 0\}$, and for \mathcal{H}^N -a.e. $(t, x) \in \Sigma$ the strong traces $\tilde{w} = \gamma \varphi_\nu(u)$, $\tilde{v} = \gamma V_{\varphi_\nu}(u)$ on Σ of the functions $\varphi_\nu(u)$, $V_{\varphi_\nu}(u)$ verify

$$(\tilde{v}(t, x), \tilde{w}(t, x)) \in \tilde{\beta} \circ V_{\varphi_\nu}^{-1}. \quad (7)$$

Notice that if V_{φ_ν} is invertible (which is the case under assumption (3)), then requirement (7) is equivalent to the requirement that $(\tilde{u}(t, x), \tilde{w}(t, x)) \in \tilde{\beta}$, where \tilde{u} is the strong trace of u on Σ , $\tilde{u} = V_{\varphi_\nu}^{-1}(\tilde{v})$.

Definition 2 can be reformulated so that to extend the entropy inequalities up to the boundaries Σ and $\{t = 0\}$. Indeed, we have

Proposition 2. *A function $u \in L^\infty(Q)$ such that strong traces $\tilde{v} := \gamma V_{\varphi_\nu}(u)$, $\tilde{w} := \gamma \varphi_\nu(u)$ on Σ exist and satisfy (7) is an entropy solution for (H) if and only if it satisfies $\forall k \in \mathbb{R}$, $\forall \xi \in C_c^\infty([0, T] \times \mathbb{R}^N)$, $\xi \geq 0$:*

³ A function $u \in L^\infty(Q)$ is called a quasi-solution if $\forall k \in \mathbb{R}$ $\eta_k^\pm(u)_t + \operatorname{div} q_k^\pm(u) = -\mu_k^\pm$ in $\mathcal{D}'(Q)$, where (η_k^\pm, q_k^\pm) are the Kruzhkov entropy-flux pairs and μ_k^\pm are Borel measures on Q , locally finite up to the boundary; see [P05, P06].

$$\begin{aligned}
0 \leq & \int_Q (u - k)^\pm \xi_t + \int_\Omega (u_0 - k)^\pm \xi(0) + \int_Q \text{sign}^\pm(u - k) f \xi \\
& + \int_Q \text{sign}^\pm(u - k) (\varphi(u) - \varphi(k)) \cdot D\xi - \int_\Sigma Q^\pm(\tilde{v}, V_{\varphi_\nu}(k)) \xi. \quad (8)
\end{aligned}$$

For the proof, one truncates $\xi \in C_c^\infty([0, T] \times \mathbb{R}^N)$ in a neighborhood of the boundaries, and passes to the limit using Definition 1.

Notice that both formulations (6),(7) and (8),(7) make sense for L^∞ (and even more general) data u_0, f , for general domain Ω with Lipschitz deformable boundary, and without assumptions (1)-(3). Uniqueness and comparison results of the next section remain valid in this general framework. For the existence part, we use a chain of approximations of (H) enjoying convenient compactness properties (this is the aim of our simplifying assumptions). We show that they converge to an entropy solution inside Q , then deduce the existence of strong traces \tilde{v}, \tilde{w} by Proposition 1, and finally, we identify the couple (\tilde{v}, \tilde{w}) as belonging to $\tilde{\beta} \circ V_{\varphi_\nu}^{-1}$.

3.1 Comparison and uniqueness of entropy solutions

Theorem 1 (The Kato inequality). *For $i = 1, 2$, let u_i be an entropy solution for Problem (H) with data $(u_0^i, f_i) \in L^1(\Omega) \times L^1(Q)$. Then for all $\xi \in C_c^\infty([0, T] \times \mathbb{R}^N), \xi \geq 0$*

$$\begin{aligned}
& \int_Q (u_1 - u_2)^+ \xi_t + \int_Q \text{sign}^+(u_1 - u_2) (\varphi(u_1) - \varphi(u_2)) \cdot D\xi \\
& + \int_Q \text{sign}^+(u_1 - u_2) (f_1 - f_2) \xi \geq \int_\Sigma (\tilde{w}_1 - \tilde{w}_2)^+ \xi. \quad (9)
\end{aligned}$$

Proof. We use the Kruzhkov method of doubling of variables. As $u_1(t, x)$, resp., $u_2(s, y)$ is an entropy solution with data $u_0^1(x)$ and $f_1(t, x)$, resp., $u_0^2(y)$ and $f_2(s, y)$, then for all $\phi = \phi(t, x, s, y) \in C_c^\infty(Q \times Q)$ one has

$$\begin{aligned}
& \int_{Q \times Q} (u_1 - u_2)^+ (\phi_t + \phi_s) + \text{sign}^+(u_1 - u_2) (\varphi(u_1) - \varphi(u_2)) \cdot (D_x \phi + D_y \phi) \\
& + \int_{Q \times Q} \text{sign}^+(u_1 - u_2) (f_1 - f_2) \geq 0. \quad (10)
\end{aligned}$$

Let $\xi \in C_c^\infty([0, T] \times \mathbb{R}^N), \xi \geq 0$ and ρ_n , resp., ρ_m , be a classical sequence of mollifiers in \mathbb{R}^N , resp., in \mathbb{R} . Define $\phi(t, x, s, y) = \mu_\delta(x) \mu_\eta(y) \rho_n(x - y) \rho_m(t - s)$. Using ϕ as a test function in (10) and passing to the limit with $\delta, \eta \rightarrow 0$ yields

$$\begin{aligned}
0 \leq & \int_{Q \times Q} \rho_m \rho_n (u_1 - u_2)^+ \xi_t + \int_{Q \times Q} \rho_m \rho_n \text{sign}^+(u_1 - u_2) (\varphi(u_1) - \varphi(u_2)) \cdot D_x \xi \\
& + \int_{Q \times Q} \text{sign}^+(u_1 - u_2) (f_1 - f_2) \rho_m \rho_n \xi \\
& - \int_{\Sigma \times Q} Q^+(\tilde{v}_1, V_{\varphi_\nu}(u_2)) \rho_m \rho_n \xi - \int_{Q \times \Sigma} Q^+(V_{\varphi_\nu}(u_1), \tilde{v}_2) \rho_m \rho_n \xi. \quad (11)
\end{aligned}$$

By Proposition 1, each of the last two terms converges to $\frac{1}{2} \int_{\Sigma} Q^+(\tilde{v}_1, \tilde{v}_2)$ as $n, m \rightarrow \infty$. Relations (7) for u_1, u_2 and the definitions of Q^+ and $\tilde{\beta}$ yield (10).

Corollary 1. *Assume (2). For data (u_0^i, f_i) , $i = 1, 2$, satisfying*

$$\begin{cases} u_0 \in (L^1 \cap L^\infty)(\Omega), f \in L^1(Q) & \text{with} \\ f(t, \cdot) \in L^\infty(\Omega) \text{ a.e. } t \in (0, T), \int_0^T \|f(t, \cdot)\|_\infty dt < \infty, \end{cases} \quad (12)$$

let u_i an entropy solution for Problem (H). Then for all $t \in (0, T)$,

$$\int_0^t \int_{\partial\Omega} (\tilde{w}_1 - \tilde{w}_2)^+ + \int_{\Omega} (u_1 - u_2)^+(t) \leq \int_{\Omega} (u_0^1 - u_0^2)^+ + \int_0^t \int_{\Omega} (f_1 - f_2)^+.$$

Thus if $u_0^1 \leq u_0^2$ a.e. on Ω and if $f_1 \leq f_2$ a.e. on Q , then $u_1 \leq u_2$ a.e. on Q . In particular, an entropy solution of (H) with data (12) is unique.

Proof. Take in (9), $\xi(t, x) = \xi_\alpha(x)\kappa(t)$, where $\kappa \in C_c^\infty([0, T])$, $\kappa \geq 0$, $\xi_\alpha \rightarrow 1$ in Ω , $|D\xi_\alpha| \leq C$ and $\text{Supp}(D\xi_\alpha) \subset \{x \mid \alpha < |x| < \alpha + 1\}$. As $\alpha \rightarrow 0$ the claim follows, because $u_i \in L^1(Q)$ under assumption (12), and φ is Lipschitz.

Notice that the comparison and uniqueness result of Corollary 1 holds true also in the L^∞ framework (see the results of Kruzhkov-Panov, Bénilan-Kruzhkov).

3.2 Existence of entropy solutions

We infer the existence of an entropy solution to Problem (H) using the tools of the nonlinear semigroup theory (cf. e.g. [BCP]). We first study the boundary problem in Ω associated with (H) (known as the “stationary” problem)

$$(S)(f) \begin{cases} u + \text{div } \varphi(u) = f & \text{on } \Omega \\ \varphi(u) \cdot \nu \in \beta(u) & \text{on } \partial\Omega. \end{cases}$$

Definition 3. *A function $u \in L^\infty(\Omega)$ is an entropy solution of Problem (S)(f) if $\forall k \in \mathbb{R}, \forall \phi \in C_c^\infty(\Omega), \phi \geq 0$ the local entropy inequalities*

$$\int_{\Omega} \text{sign}^\pm(u - k)(\varphi(u) - \varphi(k)) \cdot D\phi + \int_{\Omega} \text{sign}^\pm(u - k)(f - u)\phi \geq 0$$

hold, and traces $\tilde{v} = \gamma V_{\varphi_\nu}(u)$, $\tilde{w} = \gamma \varphi_\nu(u)$ verify (7) for \mathcal{H}^{N-1} - a.e. $x \in \partial\Omega$.

Notice that Definition 3 can be reformulated in the way of Proposition 2. Define the operator \mathcal{A} associated with the problem (S)(f) by its graph :

$$(u, f) \in \mathcal{A} \Leftrightarrow u \text{ is an entropy solution of } (S)(f + u).$$

Theorem 2. *Let (1), (2) and (3) hold. Then the operator \mathcal{A} is T -accretive with dense domain in $L^1(\Omega)$, and we have $(L^1 \cap L^\infty)(\Omega) \subset \text{Range}(I + \mathcal{A})$. Moreover, for all $(u_i, f_i) \in \mathcal{A}$, $i = 1, 2$ we have*

$$\int_{\partial\Omega} (\tilde{w}_1 - \tilde{w}_2)^\pm + \int_{\Omega} (u_1 - u_2)^\pm \leq \int_{\Omega} \text{sign}^\pm(u_1 - u_2)(f_1 - f_2) + \int_{[u_1 = u_2]} (f_1 - f_2)^\pm. \quad (13)$$

A proof of this theorem is given in [S] (see also [AS]).

Proof: The T -accretivity and (13) are obtained in the same way as Corollary 1. Let us sketch the proof of existence with data $f \in (L^1 \cap L^\infty)(\Omega)$. We use the standard vanishing viscosity approximation of the equation in $(S)(f)$ together with a Lipschitz regularization β^ε of the graph β . Let $(u^\varepsilon)_\varepsilon$ denote the corresponding sequence of approximate solutions; $\beta^\varepsilon(u^\varepsilon)$ are the corresponding normal fluxes on the boundary. Assumption (1) yields a uniform L^∞ bound on u^ε . Assumption (3) is sufficient for the strong precompactness of $(u^\varepsilon)_\varepsilon$ in $L^1_{loc}(\Omega)$, by the result of Panov [P94] (see also the well-known result of Lions-Perthame-Tadmor). This is sufficient to get (6) and deduce (by the ‘‘stationary’’ analogue of Proposition 1) the existence of boundary traces \tilde{v}, \tilde{w} of $V_{\varphi_\nu}(u), \varphi_\nu(u)$, respectively, where u is an accumulation point of $(u^\varepsilon)_\varepsilon$. It remains to show (7). Due to the flatness of $\partial\Omega$, problem $(S)(f)$ is invariant (taking into account translations of f) with respect to translations in directions $(0, h')$, $h' \in \mathbb{R}^{N-1}$. Using the analogue of (13) which holds true for $\varepsilon > 0$, by the Fréchet-Kolmogorov theorem we deduce that the sequence $(\beta^\varepsilon(u^\varepsilon))_\varepsilon$ is strongly compact in $L^1(\partial\Omega)$. We then deduce (7) by using the argument sketched in the proof of Lemma 3 below (with $\tilde{\beta}$ replaced by $\tilde{\beta}$, and $\tilde{\beta}$ replaced by β). Finally, $\overline{D(\mathcal{A})} = L^1(\Omega)$, since we show that the solution u_α to $u + \alpha \mathcal{A}(u) = f \in C_c^\infty(\Omega)$ converges to f in $L^1(\Omega)$ as $\alpha \rightarrow 0$.

Theorem 2 means that the closure of \mathcal{A} is an m - T -accretive densely defined operator in $L^1(\Omega)$ (see e.g. [BCP]). By the Crandall-Liggett theorem, it generates a T -contractive semigroup on $L^1(\Omega)$; more exactly, we have

Theorem 3. *Let (1), (2) and (3) hold. Then for any $f \in L^1(Q)$, $u_0 \in L^1(\Omega)$ there exists a unique mild solution of the abstract Cauchy problem*

$$u_t + \mathcal{A}u \ni f, \quad u(0) = u_0. \quad (14)$$

We deduce existence for Problem (H) by showing that any mild solution of problem (14) is also a solution of (H) in the sense of Definition 2.

Theorem 4. *Assume (1), (2), (3). If data (u_0, f) satisfy (12), then the mild solution of the Cauchy problem (14) is an entropy solution of Problem (H).*

Proof. For $m \in \mathbb{N}^*$, set $\varepsilon = \frac{T}{m}$. For $i = 0, 1, \dots, m$, set $t_i = \varepsilon i$ and $f_i^\varepsilon(x) = \frac{1}{\varepsilon} \int_{t_{i-1}}^{t_i} f(t, x) dt$ a.e. $x \in \Omega$. We have $\varepsilon \sum_{i=1}^m \|f_i^\varepsilon\|_{L^\infty(\Omega)} \leq \int_0^T \|f(t)\|_{L^\infty(\Omega)} dt$ and $\sum_{i=1}^m \int_{t_{i-1}}^{t_i} \|f(t) - f_i^\varepsilon\|_{L^1(\Omega)} \rightarrow 0$ as $\varepsilon \rightarrow 0$. Take $u_0^\varepsilon \in D(\mathcal{A})$ such that $\|u_0^\varepsilon - u_0\|_{L^1} \leq \varepsilon$. Let u_i^ε be the solution of $\varepsilon f_i^\varepsilon + u_{i-1}^\varepsilon \in (I + \varepsilon \mathcal{A})(u_i^\varepsilon)$,

$i = 1, \dots, m$. Set $u^\varepsilon(t) = u_i^\varepsilon, f^\varepsilon(t) = f_i^\varepsilon, \tilde{v}^\varepsilon(t) = \tilde{v}_i^\varepsilon, \tilde{w}^\varepsilon(t) = \tilde{w}_i^\varepsilon$ if $t_{i-1} \leq t \leq t_i$ ($1 \leq i \leq m$). By the nonlinear semigroup theory (see e.g. [BCP]), the sequence $(u_\varepsilon)_\varepsilon$ is precompact in $L^\infty(0, T; L^1(\Omega))$. Hence there exist a subsequence, still denoted by u_ε , and a function $u \in C(0, T; L^1(\Omega))$ such that $\|u^\varepsilon - u\|_{L^\infty(0, T; L^1(\Omega))} \rightarrow 0$ as $\varepsilon \rightarrow 0$. By construction, u admits u_0 for the strong trace on $\{t = 0\}$. Thanks to (12) and Assumption (1), $(u^\varepsilon)_\varepsilon$ is bounded in $L^\infty(Q)$. The continuation of the proof is divided into three lemmas.

Lemma 1. *The function u verifies inequality (6).*

Proof. By Definition 3, for all $\phi \in C_c^\infty(Q)$, $\phi \geq 0$, $k \in \mathbb{R}$, u_ε satisfies

$$\begin{aligned} 0 \leq & \int_Q \left(\frac{1}{\varepsilon} (u^\varepsilon(t - \varepsilon) - u^\varepsilon(t)) + f^\varepsilon(t) \right) \text{sign}^\pm(u^\varepsilon(t) - k) \phi(t) \\ & + \int_Q \text{sign}^\pm(u^\varepsilon(t) - k) (\varphi(u^\varepsilon(t)) - \varphi(k)) \cdot D\phi(t). \end{aligned} \quad (15)$$

As $\varepsilon \rightarrow 0$ the result follows, because $\text{sign}^\pm(\cdot - k)$ belongs to $\partial(\cdot - k)^\pm$:

$$\text{sign}^\pm(u^\varepsilon(t) - k) (u^\varepsilon(t - \varepsilon) - u^\varepsilon(t)) \leq (u^\varepsilon(t - \varepsilon) - k)^\pm - (u^\varepsilon(t) - k)^\pm.$$

Lemma 2. *The sequence $(\tilde{w}^\varepsilon(t, x))_\varepsilon$ of strong traces of $\varphi_\nu(u^\varepsilon)$ on $\underline{\Sigma}$ converges (up to a subsequence) \mathcal{H}^N -a.e. on Σ and in $L^1_{loc}(\Sigma)$ to $b(t, x) \in \text{Range}(\tilde{\beta})$.*

Proof. We prove that the sequence $(\tilde{w}^\varepsilon)_\varepsilon$ is bounded in $L^1(\Sigma)$, and

$$\int_0^{T-\Delta t} \|\tilde{w}^\varepsilon(t + \Delta t) - \tilde{w}^\varepsilon(t)\|_{L^1(\partial\Omega)} \leq \tilde{\psi}(\Delta t), \quad \lim_{h \rightarrow 0} \tilde{\psi}(h) = 0.$$

A similar estimate of the space translates of $\tilde{w}^\varepsilon(t, x')$ follows from (13) and the translation invariance of $\partial\Omega$. We then apply the Fréchet-Kolmogorov theorem.

Lemma 3. *Strong traces $\tilde{v} = \gamma V_{\varphi_\nu}(u)$, $\tilde{w} = \gamma \varphi_\nu(u)$ on Σ of functions $V_{\varphi_\nu}(u)$ and $\varphi_\nu(u)$ exist. Moreover, $(\tilde{v}(t, x), \tilde{w}(t, x)) \in \tilde{\beta} \circ V_{\varphi_\nu}^{-1} \mathcal{H}^N$ -a.e. $(t, x) \in \Sigma$.*

Proof. By Proposition 1, strong traces of functions $V_{\varphi_\nu}(u)$ and $\varphi_\nu(u)$ exist. Now let us prove that $(\tilde{v}, \tilde{w}) \in \tilde{\beta} \circ V_{\varphi_\nu}^{-1}$ a.e. on Σ . By choosing in inequality (3) the test function $\phi(1 - \mu_\delta)$ with $\phi \in C_c^\infty([0, T] \times \mathbb{R}^N)$, $\phi \geq 0$, $(\mu_\delta)_\delta \in C^2(\bar{\Omega})$, $\mu_\delta \rightarrow 1$ on Ω , $\mu_\delta = 0$ on $\partial\Omega$, and by letting $\varepsilon \rightarrow 0$ then $\delta \rightarrow 0$ we get

$$\int_\Sigma Q^\pm(\tilde{v}, V_{\varphi_\nu}(k)) \phi \geq \int_\Sigma (b - \varphi \circ V_{\varphi_\nu}^{-1}(V_{\varphi_\nu}(k))) \theta_k^\pm \phi, \quad (16)$$

where $\theta_k^\pm, k \in \mathbb{Q}$ is the weak-star limit in $L^\infty(\Sigma)$ of $\text{sign}^\pm(u_\varepsilon - k)$. But for a.e. $(t, x) \in \Sigma$, we have $b(t, x) = \lim_{\varepsilon \rightarrow 0} \tilde{w}^\varepsilon(t, x)$. By construction, $(\tilde{v}^\varepsilon, \tilde{w}^\varepsilon) \in \tilde{\beta} \circ V_{\varphi_\nu}^{-1}$; therefore for all $k \notin \tilde{\beta}^{-1}(b(t, x))$ we can identify $\theta_k^\pm(t, x)$ with $\text{sign}^\pm(b(t, x) - \tilde{\beta}(k))$. For simplicity, let us use assumption (3) (notice that Lemma 3 holds without (3)). We can use $\tilde{u} = V_{\varphi_\nu}^{-1}(\tilde{v})$; then (16) implies $\text{sign}^\pm(\tilde{u} - k)(\tilde{w} - \varphi_\nu(k)) \geq \text{sign}^\pm(b - k)(b - \varphi_\nu(k)) \forall k \notin \tilde{\beta}^{-1}(b)$, \mathcal{H}^N -a.e. on Σ .

Considering different values of k , we get $(\tilde{u}, \tilde{w}) \in \tilde{\beta}$. Since $\tilde{\beta} = \tilde{\beta}$, (7) holds.

Conclusion and generalizations

We justify that the entropy formulation (6),(7) or, equivalently, (8),(7) is an adequate interpretation of the formal boundary value problem (H). Uniqueness, comparison and L^1 contraction properties hold for this formulation.

Furthermore, the proof of Lemma 3 actually shows how the “effective BC” graph $\tilde{\beta}$ appears from the “formal BC” graph β (or β^ε , if the graphs are perturbed). In our proof, this passage requires strong L^1 compactness of the sequence $(\beta^\varepsilon(u^\varepsilon))_\varepsilon$ of the associated boundary fluxes.

Using the same techniques and the fact that $\tilde{\beta} = \tilde{\beta}$, one deduces existence of a solution verifying (6),(7) under weaker assumptions. For instance, if we approximate $u_0 \in L^\infty(\Omega)$ by $u_0^{m,n} = u_0^+ \mathbb{1}_{\{\|x\| < n\}} + u_0^- \mathbb{1}_{\{\|x\| < m\}} \in (L^1 \cap L^\infty)(\Omega)$, the sequence of the resulting solutions $u^{m,n}$ of (H) is monotone in each of the indices m, n , by (13). Thus we have convergence of (a subsequence of) $u^{m,n}$ in $L^1_{loc}(Q)$, of $\tilde{\beta}(u^{m,n})$ in $L^1_{loc}(\Sigma)$. Whence existence for (H) in the L^∞ framework follows (cf. [S]). Similarly, perturbation of φ by piecewise strictly monotone φ^ε permits to bypass (3), by using the techniques of L^1 continuous dependence of entropy solutions on the flux function (see e.g. the papers by Bouchut-Perthame and Chen-Karlsen). Further, approximation of β by quickly growing at infinity graphs β^ε such that $|\beta^\varepsilon| \downarrow |\beta|$ permits to get rid of hypothesis (1) and justify the remark on the extension of β in Example 1(ii). Fine translation techniques permit to consider inhomogeneous and mixed BC. Finally, in order to extend the well-posedness results to the L^1 data, one has to apply the same techniques to renormalized solutions of (H) (see Bénilan-Carrillo-Wittbold) instead of the entropy solutions.

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