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# Strong boundary traces and well-posedness for scalar conservation laws with dissipative boundary conditions

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## 1 Introduction

The aim of this paper is to give sense to the following formal problem for a scalar conservation law with boundary condition (BC, in the sequel) :

$$(H) \begin{cases} u_t + \operatorname{div} \varphi(u) = f & \text{in } Q := (0, T) \times \Omega \\ u(0, \cdot) = u_0 & \text{on } \Omega \\ \varphi_\nu(u) := \varphi(u) \cdot \nu \in \beta(u) & \text{on } \Sigma := (0, T) \times \partial\Omega. \end{cases}$$

Here  $\Omega = \mathbb{R}^+ \times \mathbb{R}^{N-1}$  ( $N \geq 1$ ),  $T > 0$ ,  $\nu$  is the unit outward normal vector on  $\partial\Omega$ ,  $\varphi : \mathbb{R} \rightarrow \mathbb{R}^N$  is continuous, and  $\beta$  is a maximal monotone graph on  $\mathbb{R}$ . Assume  $\varphi_\nu(0) - \beta(0) \ni 0$  and normalize  $\varphi, \beta$  by  $\varphi(0) = 0$ ,  $0 \in \beta(0)$ . The classical Neumann (zero-flux) and Dirichlet homogeneous BC correspond to the graphs  $\beta = \mathbb{R} \times \{0\}$  and  $\beta = \{0\} \times \mathbb{R}$ , respectively.

To show existence, we restrict our attention to the case of  $\Omega$  with flat boundary, to  $L^1 \cap L^\infty$  data  $u_0, f$  (see (12) for the precise assumption on the data), and make the following simplifying assumptions:

$$\text{there exists a constant } C \text{ such that } |\beta(z)| \geq \operatorname{sign}(z) \varphi_\nu(z) \quad \forall |z| > C; \quad (1)$$

$$\varphi \text{ is Lipschitz continuous, and } \varphi_\nu = \varphi \cdot \nu \text{ is piecewise monotone;} \quad (2)$$

$$\forall \xi \in \mathbb{R}^N \setminus \{0\}, \text{ the function } z \mapsto \xi \cdot \varphi(z) \text{ is non-constant on any interval.} \quad (3)$$

In the Conclusion, possible generalizations are indicated.

It is well known that existence for (H) generally fails if one interprets the Dirichlet BC literally (cf. [BLN]). This is also the case for general  $\beta$ , except for some particular situations (cf. e.g. [BFK]). If one approximates (H) by a sequence of problems  $(H^\epsilon)$  with the BC understood literally (e.g. parabolic “viscous” approximations, or numerical schemes), a boundary layer can form

in the corresponding solutions  $u^\epsilon$ . The convergence of  $u^\epsilon$  then takes place only locally inside the domain, and the limiting function  $u$  satisfies to the scalar conservation law with a different BC, which we call “effective BC”.

It is the goal of this paper to investigate the form of the effective BC for a general BC given by a maximal monotone graph  $\beta$ . Note that the monotonicity is necessary, if we hope the boundary condition to be dissipative in  $L^1$  (in particular, if we hope for an  $L^1$  contraction principle for solutions of (H), as in the classical situations of  $\Omega = \mathbb{R}^N$ , or of bounded  $\Omega$  with Dirichlet BC).

We suggest that the effective BC is given by a monotone graph  $\tilde{\beta}$  defined by

$$\tilde{\beta} := \left\{ \begin{array}{l} \exists b \in \overline{\text{Range}(\beta)} \text{ such that } \varphi_\nu(z) = b \text{ and} \\ (z, \varphi_\nu(z)); \text{ if } z < m := \inf \beta^{-1}(b), \text{ then } \varphi_\nu(k) \geq b \ \forall k \in [z, m[ \\ \text{if } z > M := \sup \beta^{-1}(b), \text{ then } \varphi_\nu(k) \leq b \ \forall k \in ]M, z] \end{array} \right\}, \quad (4)$$

which can be visualized as the horizontal projection of  $\beta$  on the graph of  $\varphi_\nu$ . The graph  $\tilde{\beta}$  is monotone, and  $\tilde{\tilde{\beta}} = \tilde{\beta}$  (thus operation  $\tilde{\cdot}$  is indeed a projection).

*Example 1.* (i) If  $\beta = \{0\} \times \mathbb{R}$ , then  $\tilde{\beta}$  is the Bardos-LeRoux-Nédélec graph:

$$\tilde{\beta} = \{(z, \varphi_\nu(z)) \mid \text{sign}(z)(\varphi_\nu(z) - \varphi_\nu(k)) \geq 0 \ \forall k \in [0 \wedge z, 0 \vee z]\}. \quad (5)$$

(ii) If  $\beta = \mathbb{R} \times \{0\}$ , then  $\tilde{\beta} = \{(z, \varphi_\nu(z)) \mid \varphi_\nu(z) = 0\}$ . Assumption (1) is restrictive in this case; a similar assumption is made in [BFK]. For the general case, one has to complete  $\beta$  to a maximal monotone graph on  $[-\infty, +\infty]$  before defining  $\tilde{\beta}$  by (4); see Conclusion and the forthcoming paper [AS].

## 2 Strong boundary traces for entropy solutions

The Bardos-LeRoux-Nédélec pointwise formulation of the Dirichlet BC, which can be expressed by means of the graph (5) was initially given for  $BV$  solutions (compare to the work of Otto in [MNR], where a weak formulation of the Dirichlet boundary condition is given, valid for any  $L^\infty$  solution). It has recently been realized that any  $L^\infty$  entropy solution actually has strong  $L^1_{loc}$  initial and boundary traces in a fairly general situation (see Panov [P05, P06] and the previous works of Chen-Rascle and Vasseur). In particular, the non-degeneracy condition (3) is not needed. The concept of strong trace (in  $L^1_{loc}$ ) is stated in the following definition. Set  $x = (x_1, x')$ ,  $x_1 \in \mathbb{R}^+$ ,  $x' \in \mathbb{R}^{N-1}$ .

**Definition 1.** A function  $\tilde{v} \in L^1_{loc}(\mathbb{R}^{N-1})$  is a strong trace of function  $v \in L^1_{loc}(\mathbb{R}^+ \times \mathbb{R}^{N-1})$  on  $\{x_1 = 0\}$  if for all  $\xi \in C_c(\mathbb{R}^{N-1})$ ,  $\xi \geq 0$

$$\text{ess-} \lim_{x_1 \rightarrow 0} \int_{\mathbb{R}^{N-1}} \xi(x') |v(x_1, x') - \tilde{v}(x')| dx' = 0.$$

In the sequel  $\gamma : L^1_{loc}(\mathbb{R}^+ \times \mathbb{R}^{N-1}) \rightarrow L^1_{loc}(\mathbb{R}^{N-1})$  is the strong trace operator in the sense of Definition 1. Traces of  $L^1_{loc}(Q)$  functions on  $\Sigma = (0, T) \times \mathbb{R}^{N-1}$  and on  $\{t = 0\}$  are defined similarly. Clearly, the strong trace operators are unbounded. The following proposition, inferred by the result of Panov [P05] on traces of entropy (quasi-)solutions at  $\{t = 0\}$ , is therefore remarkable.

**Proposition 1.** *Assume (2). Let  $u$  be a quasi-solution for  $u_t + \operatorname{div} \varphi(u)^3$ . Then there exists a strong trace  $\gamma V_{\varphi_\nu}(u)$  of the function  $V_{\varphi_\nu}(u)$  on  $\Sigma$  in the sense of Definition 1, where  $V_{\varphi_\nu}(\cdot)$  is the variation function of  $\varphi_\nu$  :*

$$V_{\varphi_\nu}(z) = \operatorname{sign}(z) \operatorname{Var}_{[0 \wedge z, 0 \vee z]} \varphi_\nu(\cdot).$$

In particular,  $\operatorname{sign}^\pm(u - k)(\varphi(u) - \varphi(k)) \cdot \nu$ , which can be rewritten as

$$Q^\pm(V_{\varphi_\nu}(u), V_{\varphi_\nu}(k)) := \operatorname{sign}^\pm(V_{\varphi_\nu}(u) - V_{\varphi_\nu}(k)) (\Psi_\nu(V_{\varphi_\nu}(u)) - \Psi_\nu(V_{\varphi_\nu}(k)))$$

with  $\Psi_\nu := \varphi_\nu \circ V_{\varphi_\nu}^{-1}$ , has the strong trace  $Q^\pm(\gamma V_{\varphi_\nu}(u), V_{\varphi_\nu}(k))$  on  $\Sigma$ .

Notice that the above result does not depend on (1),(3). It still holds if  $\Omega$  is a domain with piecewise  $C^1$ -smooth boundary and  $\varphi \in BV$  (see Panov [P06]).

### 3 Entropy solutions and well-posedness

**Definition 2.** *A function  $u \in L^\infty(Q)$  is said an entropy solution for Problem (H) if  $\forall k \in \mathbb{R}$ ,  $\forall \xi \in C_c^\infty(Q)$ ,  $\xi \geq 0$  the local entropy inequalities hold :*

$$\int_Q (u - k)^\pm \xi_t + \int_Q \operatorname{sign}^\pm(u - k)(\varphi(u) - \varphi(k)) \cdot D\xi + \int_Q f \operatorname{sign}^\pm(u - k) \xi \geq 0, \quad (6)$$

*u has the strong trace  $u_0$  on  $\{t = 0\}$ , and for  $\mathcal{H}^N$ -a.e.  $(t, x) \in \Sigma$  the strong traces  $\tilde{w} = \gamma \varphi_\nu(u)$ ,  $\tilde{v} = \gamma V_{\varphi_\nu}(u)$  on  $\Sigma$  of the functions  $\varphi_\nu(u)$ ,  $V_{\varphi_\nu}(u)$  verify*

$$(\tilde{v}(t, x), \tilde{w}(t, x)) \in \tilde{\beta} \circ V_{\varphi_\nu}^{-1}. \quad (7)$$

Notice that if  $V_{\varphi_\nu}$  is invertible (which is the case under assumption (3)), then requirement (7) is equivalent to the requirement that  $(\tilde{u}(t, x), \tilde{w}(t, x)) \in \tilde{\beta}$ , where  $\tilde{u}$  is the strong trace of  $u$  on  $\Sigma$ ,  $\tilde{u} = V_{\varphi_\nu}^{-1}(\tilde{v})$ .

Definition 2 can be reformulated so that to extend the entropy inequalities up to the boundaries  $\Sigma$  and  $\{t = 0\}$ . Indeed, we have

**Proposition 2.** *A function  $u \in L^\infty(Q)$  such that strong traces  $\tilde{v} := \gamma V_{\varphi_\nu}(u)$ ,  $\tilde{w} := \gamma \varphi_\nu(u)$  on  $\Sigma$  exist and satisfy (7) is an entropy solution for (H) if and only if it satisfies  $\forall k \in \mathbb{R}$ ,  $\forall \xi \in C_c^\infty([0, T] \times \mathbb{R}^N)$ ,  $\xi \geq 0$  :*

<sup>3</sup> A function  $u \in L^\infty(Q)$  is called a quasi-solution if  $\forall k \in \mathbb{R}$   $\eta_k^\pm(u)_t + \operatorname{div} q_k^\pm(u) = -\mu_k^\pm$  in  $\mathcal{D}'(Q)$ , where  $(\eta_k^\pm, q_k^\pm)$  are the Kruzhkov entropy-flux pairs and  $\mu_k^\pm$  are Borel measures on  $Q$ , locally finite up to the boundary; see [P05, P06].

$$\begin{aligned}
0 \leq & \int_Q (u - k)^\pm \xi_t + \int_\Omega (u_0 - k)^\pm \xi(0) + \int_Q \text{sign}^\pm(u - k) f \xi \\
& + \int_Q \text{sign}^\pm(u - k) (\varphi(u) - \varphi(k)) \cdot D\xi - \int_\Sigma Q^\pm(\tilde{v}, V_{\varphi_\nu}(k)) \xi. \quad (8)
\end{aligned}$$

For the proof, one truncates  $\xi \in C_c^\infty([0, T] \times \mathbb{R}^N)$  in a neighborhood of the boundaries, and passes to the limit using Definition 1.

Notice that both formulations (6),(7) and (8),(7) make sense for  $L^\infty$  (and even more general) data  $u_0, f$ , for general domain  $\Omega$  with Lipschitz deformable boundary, and without assumptions (1)-(3). Uniqueness and comparison results of the next section remain valid in this general framework. For the existence part, we use a chain of approximations of (H) enjoying convenient compactness properties (this is the aim of our simplifying assumptions). We show that they converge to an entropy solution inside  $Q$ , then deduce the existence of strong traces  $\tilde{v}, \tilde{w}$  by Proposition 1, and finally, we identify the couple  $(\tilde{v}, \tilde{w})$  as belonging to  $\tilde{\beta} \circ V_{\varphi_\nu}^{-1}$ .

### 3.1 Comparison and uniqueness of entropy solutions

**Theorem 1 (The Kato inequality).** *For  $i = 1, 2$ , let  $u_i$  be an entropy solution for Problem (H) with data  $(u_0^i, f_i) \in L^1(\Omega) \times L^1(Q)$ . Then for all  $\xi \in C_c^\infty([0, T] \times \mathbb{R}^N), \xi \geq 0$*

$$\begin{aligned}
& \int_Q (u_1 - u_2)^+ \xi_t + \int_Q \text{sign}^+(u_1 - u_2) (\varphi(u_1) - \varphi(u_2)) \cdot D\xi \\
& + \int_Q \text{sign}^+(u_1 - u_2) (f_1 - f_2) \xi \geq \int_\Sigma (\tilde{w}_1 - \tilde{w}_2)^+ \xi. \quad (9)
\end{aligned}$$

*Proof.* We use the Kruzhkov method of doubling of variables. As  $u_1(t, x)$ , resp.,  $u_2(s, y)$  is an entropy solution with data  $u_0^1(x)$  and  $f_1(t, x)$ , resp.,  $u_0^2(y)$  and  $f_2(s, y)$ , then for all  $\phi = \phi(t, x, s, y) \in C_c^\infty(Q \times Q)$  one has

$$\begin{aligned}
& \int_{Q \times Q} (u_1 - u_2)^+ (\phi_t + \phi_s) + \text{sign}^+(u_1 - u_2) (\varphi(u_1) - \varphi(u_2)) \cdot (D_x \phi + D_y \phi) \\
& + \int_{Q \times Q} \text{sign}^+(u_1 - u_2) (f_1 - f_2) \geq 0. \quad (10)
\end{aligned}$$

Let  $\xi \in C_c^\infty([0, T] \times \mathbb{R}^N), \xi \geq 0$  and  $\rho_n$ , resp.,  $\rho_m$ , be a classical sequence of mollifiers in  $\mathbb{R}^N$ , resp., in  $\mathbb{R}$ . Define  $\phi(t, x, s, y) = \mu_\delta(x) \mu_\eta(y) \rho_n(x - y) \rho_m(t - s)$ . Using  $\phi$  as a test function in (10) and passing to the limit with  $\delta, \eta \rightarrow 0$  yields

$$\begin{aligned}
0 \leq & \int_{Q \times Q} \rho_m \rho_n (u_1 - u_2)^+ \xi_t + \int_{Q \times Q} \rho_m \rho_n \text{sign}^+(u_1 - u_2) (\varphi(u_1) - \varphi(u_2)) \cdot D_x \xi \\
& + \int_{Q \times Q} \text{sign}^+(u_1 - u_2) (f_1 - f_2) \rho_m \rho_n \xi \\
& - \int_{\Sigma \times Q} Q^+(\tilde{v}_1, V_{\varphi_\nu}(u_2)) \rho_m \rho_n \xi - \int_{Q \times \Sigma} Q^+(V_{\varphi_\nu}(u_1), \tilde{v}_2) \rho_m \rho_n \xi. \quad (11)
\end{aligned}$$

By Proposition 1, each of the last two terms converges to  $\frac{1}{2} \int_{\Sigma} Q^+(\tilde{v}_1, \tilde{v}_2)$  as  $n, m \rightarrow \infty$ . Relations (7) for  $u_1, u_2$  and the definitions of  $Q^+$  and  $\tilde{\beta}$  yield (10).

**Corollary 1.** *Assume (2). For data  $(u_0^i, f_i)$ ,  $i = 1, 2$ , satisfying*

$$\begin{cases} u_0 \in (L^1 \cap L^\infty)(\Omega), f \in L^1(Q) & \text{with} \\ f(t, \cdot) \in L^\infty(\Omega) \text{ a.e. } t \in (0, T), \int_0^T \|f(t, \cdot)\|_\infty dt < \infty, \end{cases} \quad (12)$$

let  $u_i$  an entropy solution for Problem (H). Then for all  $t \in (0, T)$ ,

$$\int_0^t \int_{\partial\Omega} (\tilde{w}_1 - \tilde{w}_2)^+ + \int_{\Omega} (u_1 - u_2)^+(t) \leq \int_{\Omega} (u_0^1 - u_0^2)^+ + \int_0^t \int_{\Omega} (f_1 - f_2)^+.$$

Thus if  $u_0^1 \leq u_0^2$  a.e. on  $\Omega$  and if  $f_1 \leq f_2$  a.e. on  $Q$ , then  $u_1 \leq u_2$  a.e. on  $Q$ . In particular, an entropy solution of (H) with data (12) is unique.

*Proof.* Take in (9),  $\xi(t, x) = \xi_\alpha(x)\kappa(t)$ , where  $\kappa \in C_c^\infty([0, T])$ ,  $\kappa \geq 0$ ,  $\xi_\alpha \rightarrow 1$  in  $\Omega$ ,  $|D\xi_\alpha| \leq C$  and  $\text{Supp}(D\xi_\alpha) \subset \{x \mid \alpha < |x| < \alpha + 1\}$ . As  $\alpha \rightarrow 0$  the claim follows, because  $u_i \in L^1(Q)$  under assumption (12), and  $\varphi$  is Lipschitz.

Notice that the comparison and uniqueness result of Corollary 1 holds true also in the  $L^\infty$  framework (see the results of Kruzhkov-Panov, Bénilan-Kruzhkov).

### 3.2 Existence of entropy solutions

We infer the existence of an entropy solution to Problem (H) using the tools of the nonlinear semigroup theory (cf. e.g. [BCP]). We first study the boundary problem in  $\Omega$  associated with (H) (known as the “stationary” problem)

$$(S)(f) \begin{cases} u + \text{div } \varphi(u) = f & \text{on } \Omega \\ \varphi(u) \cdot \nu \in \beta(u) & \text{on } \partial\Omega. \end{cases}$$

**Definition 3.** *A function  $u \in L^\infty(\Omega)$  is an entropy solution of Problem (S)(f) if  $\forall k \in \mathbb{R}, \forall \phi \in C_c^\infty(\Omega), \phi \geq 0$  the local entropy inequalities*

$$\int_{\Omega} \text{sign}^\pm(u - k)(\varphi(u) - \varphi(k)) \cdot D\phi + \int_{\Omega} \text{sign}^\pm(u - k)(f - u)\phi \geq 0$$

hold, and traces  $\tilde{v} = \gamma V_{\varphi_\nu}(u)$ ,  $\tilde{w} = \gamma \varphi_\nu(u)$  verify (7) for  $\mathcal{H}^{N-1}$  - a.e.  $x \in \partial\Omega$ .

Notice that Definition 3 can be reformulated in the way of Proposition 2. Define the operator  $\mathcal{A}$  associated with the problem (S)(f) by its graph :

$$(u, f) \in \mathcal{A} \Leftrightarrow u \text{ is an entropy solution of } (S)(f + u).$$

**Theorem 2.** *Let (1), (2) and (3) hold. Then the operator  $\mathcal{A}$  is  $T$ -accretive with dense domain in  $L^1(\Omega)$ , and we have  $(L^1 \cap L^\infty)(\Omega) \subset \text{Range}(I + \mathcal{A})$ . Moreover, for all  $(u_i, f_i) \in \mathcal{A}$ ,  $i = 1, 2$  we have*

$$\int_{\partial\Omega} (\tilde{w}_1 - \tilde{w}_2)^\pm + \int_{\Omega} (u_1 - u_2)^\pm \leq \int_{\Omega} \text{sign}^\pm(u_1 - u_2)(f_1 - f_2) + \int_{[u_1 = u_2]} (f_1 - f_2)^\pm. \quad (13)$$

A proof of this theorem is given in [S] (see also [AS]).

*Proof:* The  $T$ -accretivity and (13) are obtained in the same way as Corollary 1. Let us sketch the proof of existence with data  $f \in (L^1 \cap L^\infty)(\Omega)$ . We use the standard vanishing viscosity approximation of the equation in  $(S)(f)$  together with a Lipschitz regularization  $\beta^\varepsilon$  of the graph  $\beta$ . Let  $(u^\varepsilon)_\varepsilon$  denote the corresponding sequence of approximate solutions;  $\beta^\varepsilon(u^\varepsilon)$  are the corresponding normal fluxes on the boundary. Assumption (1) yields a uniform  $L^\infty$  bound on  $u^\varepsilon$ . Assumption (3) is sufficient for the strong precompactness of  $(u^\varepsilon)_\varepsilon$  in  $L^1_{loc}(\Omega)$ , by the result of Panov [P94] (see also the well-known result of Lions-Perthame-Tadmor). This is sufficient to get (6) and deduce (by the ‘‘stationary’’ analogue of Proposition 1) the existence of boundary traces  $\tilde{v}, \tilde{w}$  of  $V_{\varphi_\nu}(u), \varphi_\nu(u)$ , respectively, where  $u$  is an accumulation point of  $(u^\varepsilon)_\varepsilon$ . It remains to show (7). Due to the flatness of  $\partial\Omega$ , problem  $(S)(f)$  is invariant (taking into account translations of  $f$ ) with respect to translations in directions  $(0, h')$ ,  $h' \in \mathbb{R}^{N-1}$ . Using the analogue of (13) which holds true for  $\varepsilon > 0$ , by the Fréchet-Kolmogorov theorem we deduce that the sequence  $(\beta^\varepsilon(u^\varepsilon))_\varepsilon$  is strongly compact in  $L^1(\partial\Omega)$ . We then deduce (7) by using the argument sketched in the proof of Lemma 3 below (with  $\tilde{\beta}$  replaced by  $\tilde{\beta}$ , and  $\tilde{\beta}$  replaced by  $\beta$ ). Finally,  $\overline{D(\mathcal{A})} = L^1(\Omega)$ , since we show that the solution  $u_\alpha$  to  $u + \alpha \mathcal{A}(u) = f \in C_c^\infty(\Omega)$  converges to  $f$  in  $L^1(\Omega)$  as  $\alpha \rightarrow 0$ .

Theorem 2 means that the closure of  $\mathcal{A}$  is an  $m$ - $T$ -accretive densely defined operator in  $L^1(\Omega)$  (see e.g. [BCP]). By the Crandall-Liggett theorem, it generates a  $T$ -contractive semigroup on  $L^1(\Omega)$ ; more exactly, we have

**Theorem 3.** *Let (1), (2) and (3) hold. Then for any  $f \in L^1(Q)$ ,  $u_0 \in L^1(\Omega)$  there exists a unique mild solution of the abstract Cauchy problem*

$$u_t + \mathcal{A}u \ni f, \quad u(0) = u_0. \quad (14)$$

We deduce existence for Problem (H) by showing that any mild solution of problem (14) is also a solution of (H) in the sense of Definition 2.

**Theorem 4.** *Assume (1), (2), (3). If data  $(u_0, f)$  satisfy (12), then the mild solution of the Cauchy problem (14) is an entropy solution of Problem (H).*

*Proof.* For  $m \in \mathbb{N}^*$ , set  $\varepsilon = \frac{T}{m}$ . For  $i = 0, 1, \dots, m$ , set  $t_i = \varepsilon i$  and  $f_i^\varepsilon(x) = \frac{1}{\varepsilon} \int_{t_{i-1}}^{t_i} f(t, x) dt$  a.e.  $x \in \Omega$ . We have  $\varepsilon \sum_{i=1}^m \|f_i^\varepsilon\|_{L^\infty(\Omega)} \leq \int_0^T \|f(t)\|_{L^\infty(\Omega)} dt$  and  $\sum_{i=1}^m \int_{t_{i-1}}^{t_i} \|f(t) - f_i^\varepsilon\|_{L^1(\Omega)} \rightarrow 0$  as  $\varepsilon \rightarrow 0$ . Take  $u_0^\varepsilon \in D(\mathcal{A})$  such that  $\|u_0^\varepsilon - u_0\|_{L^1} \leq \varepsilon$ . Let  $u_i^\varepsilon$  be the solution of  $\varepsilon f_i^\varepsilon + u_{i-1}^\varepsilon \in (I + \varepsilon \mathcal{A})(u_i^\varepsilon)$ ,

$i = 1, \dots, m$ . Set  $u^\varepsilon(t) = u_i^\varepsilon, f^\varepsilon(t) = f_i^\varepsilon, \tilde{v}^\varepsilon(t) = \tilde{v}_i^\varepsilon, \tilde{w}^\varepsilon(t) = \tilde{w}_i^\varepsilon$  if  $t_{i-1} \leq t \leq t_i$  ( $1 \leq i \leq m$ ). By the nonlinear semigroup theory (see e.g. [BCP]), the sequence  $(u_\varepsilon)_\varepsilon$  is precompact in  $L^\infty(0, T; L^1(\Omega))$ . Hence there exist a subsequence, still denoted by  $u_\varepsilon$ , and a function  $u \in C(0, T; L^1(\Omega))$  such that  $\|u^\varepsilon - u\|_{L^\infty(0, T; L^1(\Omega))} \rightarrow 0$  as  $\varepsilon \rightarrow 0$ . By construction,  $u$  admits  $u_0$  for the strong trace on  $\{t = 0\}$ . Thanks to (12) and Assumption (1),  $(u^\varepsilon)_\varepsilon$  is bounded in  $L^\infty(Q)$ . The continuation of the proof is divided into three lemmas.

**Lemma 1.** *The function  $u$  verifies inequality (6).*

*Proof.* By Definition 3, for all  $\phi \in C_c^\infty(Q)$ ,  $\phi \geq 0$ ,  $k \in \mathbb{R}$ ,  $u_\varepsilon$  satisfies

$$\begin{aligned} 0 \leq & \int_Q \left( \frac{1}{\varepsilon} (u^\varepsilon(t - \varepsilon) - u^\varepsilon(t)) + f^\varepsilon(t) \right) \text{sign}^\pm(u^\varepsilon(t) - k) \phi(t) \\ & + \int_Q \text{sign}^\pm(u^\varepsilon(t) - k) (\varphi(u^\varepsilon(t)) - \varphi(k)) \cdot D\phi(t). \end{aligned} \quad (15)$$

As  $\varepsilon \rightarrow 0$  the result follows, because  $\text{sign}^\pm(\cdot - k)$  belongs to  $\partial(\cdot - k)^\pm$ :

$$\text{sign}^\pm(u^\varepsilon(t) - k) (u^\varepsilon(t - \varepsilon) - u^\varepsilon(t)) \leq (u^\varepsilon(t - \varepsilon) - k)^\pm - (u^\varepsilon(t) - k)^\pm.$$

**Lemma 2.** *The sequence  $(\tilde{w}^\varepsilon(t, x))_\varepsilon$  of strong traces of  $\varphi_\nu(u^\varepsilon)$  on  $\underline{\Sigma}$  converges (up to a subsequence)  $\mathcal{H}^N$ -a.e. on  $\Sigma$  and in  $L^1_{loc}(\Sigma)$  to  $b(t, x) \in \text{Range}(\tilde{\beta})$ .*

*Proof.* We prove that the sequence  $(\tilde{w}^\varepsilon)_\varepsilon$  is bounded in  $L^1(\Sigma)$ , and

$$\int_0^{T-\Delta t} \|\tilde{w}^\varepsilon(t + \Delta t) - \tilde{w}^\varepsilon(t)\|_{L^1(\partial\Omega)} \leq \tilde{\psi}(\Delta t), \quad \lim_{h \rightarrow 0} \tilde{\psi}(h) = 0.$$

A similar estimate of the space translates of  $\tilde{w}^\varepsilon(t, x')$  follows from (13) and the translation invariance of  $\partial\Omega$ . We then apply the Fréchet-Kolmogorov theorem.

**Lemma 3.** *Strong traces  $\tilde{v} = \gamma V_{\varphi_\nu}(u)$ ,  $\tilde{w} = \gamma \varphi_\nu(u)$  on  $\Sigma$  of functions  $V_{\varphi_\nu}(u)$  and  $\varphi_\nu(u)$  exist. Moreover,  $(\tilde{v}(t, x), \tilde{w}(t, x)) \in \tilde{\beta} \circ V_{\varphi_\nu}^{-1} \mathcal{H}^N$ -a.e.  $(t, x) \in \Sigma$ .*

*Proof.* By Proposition 1, strong traces of functions  $V_{\varphi_\nu}(u)$  and  $\varphi_\nu(u)$  exist. Now let us prove that  $(\tilde{v}, \tilde{w}) \in \tilde{\beta} \circ V_{\varphi_\nu}^{-1}$  a.e. on  $\Sigma$ . By choosing in inequality (3) the test function  $\phi(1 - \mu_\delta)$  with  $\phi \in C_c^\infty([0, T] \times \mathbb{R}^N)$ ,  $\phi \geq 0$ ,  $(\mu_\delta)_\delta \in C^2(\bar{\Omega})$ ,  $\mu_\delta \rightarrow 1$  on  $\Omega$ ,  $\mu_\delta = 0$  on  $\partial\Omega$ , and by letting  $\varepsilon \rightarrow 0$  then  $\delta \rightarrow 0$  we get

$$\int_\Sigma Q^\pm(\tilde{v}, V_{\varphi_\nu}(k)) \phi \geq \int_\Sigma (b - \varphi \circ V_{\varphi_\nu}^{-1}(V_{\varphi_\nu}(k))) \theta_k^\pm \phi, \quad (16)$$

where  $\theta_k^\pm, k \in \mathbb{Q}$  is the weak-star limit in  $L^\infty(\Sigma)$  of  $\text{sign}^\pm(u_\varepsilon - k)$ . But for a.e.  $(t, x) \in \Sigma$ , we have  $b(t, x) = \lim_{\varepsilon \rightarrow 0} \tilde{w}^\varepsilon(t, x)$ . By construction,  $(\tilde{v}^\varepsilon, \tilde{w}^\varepsilon) \in \tilde{\beta} \circ V_{\varphi_\nu}^{-1}$ ; therefore for all  $k \notin \tilde{\beta}^{-1}(b(t, x))$  we can identify  $\theta_k^\pm(t, x)$  with  $\text{sign}^\pm(b(t, x) - \tilde{\beta}(k))$ . For simplicity, let us use assumption (3) (notice that Lemma 3 holds without (3)). We can use  $\tilde{u} = V_{\varphi_\nu}^{-1}(\tilde{v})$ ; then (16) implies  $\text{sign}^\pm(\tilde{u} - k)(\tilde{w} - \varphi_\nu(k)) \geq \text{sign}^\pm(b - k)(b - \varphi_\nu(k)) \forall k \notin \tilde{\beta}^{-1}(b)$ ,  $\mathcal{H}^N$ -a.e. on  $\Sigma$ .

Considering different values of  $k$ , we get  $(\tilde{u}, \tilde{w}) \in \tilde{\beta}$ . Since  $\tilde{\beta} = \tilde{\beta}$ , (7) holds.

## Conclusion and generalizations

We justify that the entropy formulation (6),(7) or, equivalently, (8),(7) is an adequate interpretation of the formal boundary value problem (H). Uniqueness, comparison and  $L^1$  contraction properties hold for this formulation.

Furthermore, the proof of Lemma 3 actually shows how the “effective BC” graph  $\tilde{\beta}$  appears from the “formal BC” graph  $\beta$  (or  $\beta^\varepsilon$ , if the graphs are perturbed). In our proof, this passage requires strong  $L^1$  compactness of the sequence  $(\beta^\varepsilon(u^\varepsilon))_\varepsilon$  of the associated boundary fluxes.

Using the same techniques and the fact that  $\tilde{\beta} = \tilde{\beta}$ , one deduces existence of a solution verifying (6),(7) under weaker assumptions. For instance, if we approximate  $u_0 \in L^\infty(\Omega)$  by  $u_0^{m,n} = u_0^+ \mathbb{1}_{\{\|x\| < n\}} + u_0^- \mathbb{1}_{\{\|x\| < m\}} \in (L^1 \cap L^\infty)(\Omega)$ , the sequence of the resulting solutions  $u^{m,n}$  of (H) is monotone in each of the indices  $m, n$ , by (13). Thus we have convergence of (a subsequence of)  $u^{m,n}$  in  $L^1_{loc}(Q)$ , of  $\tilde{\beta}(u^{m,n})$  in  $L^1_{loc}(\Sigma)$ . Whence existence for (H) in the  $L^\infty$  framework follows (cf. [S]). Similarly, perturbation of  $\varphi$  by piecewise strictly monotone  $\varphi^\varepsilon$  permits to bypass (3), by using the techniques of  $L^1$  continuous dependence of entropy solutions on the flux function (see e.g. the papers by Bouchut-Perthame and Chen-Karlsen). Further, approximation of  $\beta$  by quickly growing at infinity graphs  $\beta^\varepsilon$  such that  $|\beta^\varepsilon| \downarrow |\beta|$  permits to get rid of hypothesis (1) and justify the remark on the extension of  $\beta$  in Example 1(ii). Fine translation techniques permit to consider inhomogeneous and mixed BC. Finally, in order to extend the well-posedness results to the  $L^1$  data, one has to apply the same techniques to renormalized solutions of (H) (see Bénilan-Carrillo-Wittbold) instead of the entropy solutions.

## References

- [AS] Andreianov, B., Sbihi, K.: On boundary conditions for scalar conservation laws. In preparation
- [BLN] Bardos, C., Le Roux, A. Y., Nédélec, J.C.: First order quasilinear equations with boundary conditions. *Comm. Partial Diff. Equ.*, **4**, 1017–1034 (1979)
- [BCP] Bénilan, Ph., Crandall, M.G., Pazy, A.: Evolution equations governed by accretive operators. Preprint book
- [BFK] Bürger, R., Frid, H., Karlsen, K.H.: On the well-posedness of entropy solutions to conservation laws with a zero-flux boundary condition. To appear in *J. Math. Anal. Appl.*
- [MNRR] Malek, J., Nečas, J., Rokyta, M., Ružička, M.: *Weak and Measure-valued Solutions to Evolutionary PDEs.*, Chapman&Hall, 1996.
- [P94] Panov, E. Yu.: On sequences of measure-valued solutions of a first-order quasilinear equation. *Mat.Sb.*, **185**,no.2, 87–106(1994)
- [P05] Panov, E. Yu.: Existence of strong traces for generalized solutions of multidimensional scalar conservation laws. *J.Hyp.Diff.Eq.*, **4**, 885–908(2005)
- [P06] Panov, E. Yu.: communication at HYP06 Lyon, paper in preparation
- [S] Sbihi, K.: Etude de quelques E.D.P. non linéaires dans  $L^1$  avec des conditions générales sur le bord. Thesis, Strasbourg University, France (2006)