Ondes non linéaires: Bifurcations et dynamique

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NONLINEAR WAVES ON WATER



Water wave [David Sanger Photography]



Solitary wave Lagoon of Molokai, Hawaii [photo : R.I. Odom]



Roll wave Channel in Lions Bay, Canada [website of N. Balmforth]



Mascaret de St Pardon Dordogne river



Tsunami in Asia



Rogue wave Chemical tanker ship Stolt Surf [photo : K. Petersen]

OTHER NONLINEAR WAVES



Kelvin-Helmholtz clouds Mount Duval, Australia [English Wikipedia : GRAHAMUK]



Morning Glory cloud near Burketown, Australia [author : Mick Petrov]



Hurricane



Fire rainbow Northern Idaho



Sound wave Bell Telephone Laboratories [book by David C. Knight]

PATTERNS IN NATURE



Sand patterns [photo : R. Niebrugge]











The Mathematics of ... Nonlinear Waves and Patterns

- observed in nature, experiments, numerical simulations
- particular solutions of PDEs or ODEs
 - well-defined temporal and spatial structure
 - e.g., traveling waves
- play a key role in the dynamics of the underlying system

The Mathematics of ... Nonlinear Waves and Patterns

Questions

- existence spatial and temporal properties
- **stability** *spatial and temporal behavior*
- interactions
- ...
- role in the dynamics of the system

The Mathematics of ... Nonlinear Waves and Patterns

Methods

- ... many different ...
- ... not enough ...

numerical

• analytical

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FIRST EXAMPLE



WATER WAVES



WATER WAVES



WATER-WAVE PROBLEM

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- gravity-capillary water waves
 - three-dimensional inviscid fluid layer
 - constant density ρ
 - gravity and surface tension
 - irrotational flow

Second Example



Defects in patterns

- dislocations
- grain boundaries

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• disclinations

[D. Boyer, J. Viñals]

Defects in Striped Patterns

• Occur in a wide range of systems

- Rayleigh-Bénard convection experiment
- crystal patterns in material science
- chemical reactions
- biology
-

THIRD EXAMPLE ...

Collaboration FEMTO-ST & LMB



- Irina Balakireva
- Yanne K. Chembo
- Aurélien Coillet

(Lm^B)

- Cyril Godey
- Mariana Haragus

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LE PROBLÈME PHYSIQUE



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Résonateur optique



- Permet de générer des horloges de ultra-haute précision.
- Applications en aérospatiale, en télécommunications (radars, GPS...).

THE EXISTENCE PROBLEM

... A Dynamical Systems Approach

• PDEs in unbounded domains



[Kirchgässner, 1982]

• x timelike coordinate

... A Dynamical Systems Approach

• PDEs in unbounded domains



[Kirchgässner, 1982]

- x timelike coordinate
- Dynamical system
 [S1]

[SPATIAL DYNAMICS]

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$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\,\mathsf{U}=\mathsf{F}\,(\mathsf{U},\mu),\quad\mathsf{U}(\mathsf{x})\in\mathcal{X}$$

- U(x) belongs to a Hilbert (Banach) space X of functions depending upon the "space" variables;
- $\mu \in \mathbb{R}^m$ parameters.

Spatial and Radial Dynamics

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Spatial and Radial Dynamics

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Spatial and Radial Dynamics



[H. & Groves, 2003]



[Kirchgässner, 1994]



[Scheel, 2003]

SPATIAL DYNAMICS

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• Dynamical system

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\,\mathsf{U}=\mathsf{F}\,(\mathsf{U},\mu),\quad\mathsf{U}(\mathsf{x})\in\mathcal{X}$$

• Nonlinear waves are found as

bounded solutions of the dynamical system

SPATIAL DYNAMICS

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Nonlinear waves are found as

bounded solutions of the dynamical system

• What determines the shape of the wave?

SHAPE OF SOLUTIONS ...





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- ... determined by
 - boundary conditions in the space variables y
 - type of the bounded solution (localized, periodic, ... in x)

FOR INSTANCE ...



SPATIAL DYNAMICS APPROACH

1 Dynamical system

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}} \, \mathsf{U} = \mathsf{F}(\mathsf{U},\mu), \quad \mathsf{U}(\mathsf{x}) \in \mathcal{X}$$

LOCAL BIFURCATIONS

1 Dynamical system

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\,\mathsf{U}=\mathsf{F}\,(\mathsf{U},\mu),\quad\mathsf{U}(\mathsf{x})\in\mathcal{X}$$

2 Bifurcation points : critical parameter values μ_*

- start with a particular solution U_{*} (often U_{*} = 0);
- determine the spectrum of $D_UF(U_*, \mu)$
- bifurcation point μ_{*} : if the spectrum of D_UF(U_{*}, μ_{*}) contains purely imaginary values



REDUCTION

1 Dynamical system

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}} \, \mathsf{U} = \mathsf{F} \, (\mathsf{U}, \mu), \quad \mathsf{U}(\mathsf{x}) \in \mathcal{X}$$

2 Bifurcation points : critical parameter values μ_*

3 Center manifold reduction

REDUCTION

1 Dynamical system

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\mathsf{U}=\mathsf{F}(\mathsf{U},\mu), \quad \mathsf{U}(\mathsf{x})\in\mathcal{X}$$

2 Bifurcation points : critical parameter values μ_{\ast}

- **3** Center manifold reduction
 - spectrum of D_UF(U_{*}, μ_{*})
 - small bounded orbits lie on a center manifold
 - finite-dimensional center manifold
 - study the (reduced) dynamics on the center manifold
 - \rightarrow reduced ODE

[Pliss, Kelley, ..., Mielke]

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REDUCED SYSTEM

1 Dynamical system

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\mathsf{U}=\mathsf{F}(\mathsf{U},\mu), \quad \mathsf{U}(\mathsf{x})\in\mathcal{X}$$

2 Bifurcation points : critical parameter values μ_*

8 Center manifold reduction : reduced system of ODEs

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\,\mathsf{v}=\mathsf{g}\,(\mathsf{v},\mu),\quad\mathsf{v}(\mathsf{x})\in\mathbb{R}^{\mathsf{d}}$$

4 Bounded orbits of the reduced system of ODEs

• e.g., use normal form theory

[Poincaré, Birkhoff, Arnold, Elphick et al., ...]

- study a truncated system
- show persistence of the truncated dynamics







FIRST EXAMPLE

WATER WAVES


WATER WAVES



WATER WAVES



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• Domain

 $\mathsf{D}_\eta = \{(\mathsf{x},\mathsf{y},\mathsf{z}) \; : \; \mathsf{x},\mathsf{z}\in\mathbb{R}, \; \mathsf{y}\in(\mathsf{0},\mathsf{h}+\eta(\mathsf{x},\mathsf{z},\mathsf{t}))\}$

• depth at rest h

EULER EQUATIONS

$$\begin{split} \phi_{xx} + \phi_{yy} + \phi_{zz} &= 0 \quad \text{for} \quad 0 < y < 1 + \eta \\ \phi_y &= 0 \quad \text{on} \quad y = 0 \\ \phi_y &= \eta_t + \eta_x + \eta_x \phi_x + \eta_z \phi_z \quad \text{on} \quad y = 1 + \eta \\ \phi_t + \phi_x + \frac{1}{2} \left(\phi_x^2 + \phi_y^2 + \phi_z^2 \right) + \alpha \eta - \beta \mathcal{K} = 0 \quad \text{on} \quad y = 1 + \eta \end{split}$$

• velocity potential ϕ ; free surface $1+\eta$

• mean curvature
$$\mathcal{K} = \left[\frac{\eta_x}{\sqrt{1+\eta_x^2+\eta_z^2}}\right]_x + \left[\frac{\eta_z}{\sqrt{1+\eta_x^2+\eta_z^2}}\right]_z$$

- parameters
 - inverse square of the Froude number

$$\alpha = \frac{\mathsf{gh}}{\mathsf{c}^2}$$

• Weber number

$$\beta = \frac{\sigma}{\rho \mathsf{hc}^2}$$

EULER EQUATIONS

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- very rich dynamics
- difficulties
 - variable domain (free surface)
 - nonlinear boundary conditions
- symmetries, Hamiltonian structure
- many particular solutions

THE SOLITARY WAVE



John Scott Russell (1808 – 1882)



- Scottish civil engineer
- naval architect
- shipbuilder
- discovery of the solitary wave

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Experimental setup – Union Canal Edinburgh



JOHN SCOTT RUSSELL (1808 – 1882)

1834 : "The happiest day of my life"

"the boat suddently stopped – not so the mass of water in the channel which it had put in motion

a large, solitary, progressive wave"

[Recherches Hydrauliques, par M. H. Darcy et M. H. Bazin, Deuxième Partie, Paris : Imprimerie Impériale, MDCCCLXV,

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p.9]



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p.9]

Bridge 11, Hermiston Walk Heriot Watt University







AIRY AND STOKES

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Sir George Biddell Airy (1801 – 1892)

- mathematician and astronomer
- Airy wave



Sir George Gabriel Stokes (1819 –1903)

- mathematician and physicist
- Stokes wave

AIRY AND STOKES



Sir George Biddell Airy (1801 – 1892)





Sir George Gabriel Stokes (1819 –1903)



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LOOKING FOR AN EXPLANATION ...

1895 : Korteweg & de Vries

$u_t + u_x + u_{xxx} + uu_x = 0$



1872, 1877 : Boussinesq



1877 : Boussinesq

MÉMOIRES

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PRÉSENTÉS PAR DIVERS SAVANTS

À L'ACADÉMIE DES SCIENCES DE L'INSTITUT DE FRANCE.

EXTRAST DES TOMES XXIII ET XXIV.

ESSAL

LA THÉORIE DES EAUX COUBANTES.



PARIS.

IMPRIMERIE NATIONALE.

N DCCC LXXVIL

J. BOUSSINESQ. proche o proche les variations de l/, c'est-à-dire les changements que prouvera le profil longitudinal de l'onde 0/.

que provinta se prom congramma de tour en estera plus qu'à évaluer la partie non permanente U de la vitesse. On a pour cela la seconde équation du problème, (x36) ou (x70 bis) [p. 300], dans laquelle $\frac{d}{dt}$ se réduit sensiblement à $\frac{du}{dt}$. Sa comparsison à (x83) permet de poser

 $\frac{d}{ds}(hU - h'\omega) = 0$,

ou bien, en multipliant par ds et intégrant de manière que $hU - h'\omega$ se réduise à HU₀ aux points que les ondes n'ont pas encore atteints,

 $hU - h'\omega = HU_o$, c'est-h-dire $(H + h')(U_o + U') - h'\omega = HU_o$.

⁰¹ Il surait été préférable d'obtenir cette équation par l'intégration directe de (281) ou avent de parler dos vitesses de propagation as A cot effet, on annait appelé ψ₁, par exemple, l'expression

 $\psi_1 = \frac{dk'}{dt} + \omega_* \frac{dk'}{ds} + \frac{\omega_*(\omega_* - U_*)}{2\omega_* - (1 + \alpha')U_*} \frac{d}{dt} \left(\frac{2 + k}{2} \frac{k'!}{H} + \frac{k'H'}{3} \frac{d^2k'}{ds'} \right),$

et l'on sursit reconnu, au moyen de (264), que l'équation (281) revient sensiblement à



suraiost donné immédiatement, par la substitution à $--\frac{dV}{dt}$ de sa valour tirée de (883 bir), l'espression (889) de so.

LATER . . . THE SOLITARY WAVE BECOMES A SOLITON

1955 : Fermi, Pasta, Ulam recurrence **1963, 1965 : Zabuski, Kruskal** 'soliton' **1967 : Gardner, Green, Kruskal, Miura** (inverse scattering transform)

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- water waves
- nonlinear optics
- nonlinear acoustics
- plasma waves
- . . .

TODAY ... MANY DIFFERENT EXISTENCE THEORIES



Today ... Many Different Existence Theories



Today ... Many Different Existence Theories





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Euler Equations : Three-dimensional Traveling Waves



[Groves, Mielke, Craig, Nicholls, H., Kirchgässner, Deng, Sun, Sandstede, looss, Plotnikov, Wahlén, ...]

SECOND EXAMPLE

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Defects in Striped Patterns



[D. Boyer, J. Viñals]

EXISTENCE OF DEFECTS

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- dislocations
- grain boundaries
- disclinations

Defects in Striped Patterns

- Occur in a wide range of systems
- Existence studies

 in the frame of modulation equations, e.g., the Newell-Whitehead-Segel equation
[Boyer, Viñals, Manneville, Pomeau, Newell, Passot, Bowman, Malomed, Nepomnyashchy, Trybelsky, Ercolani, Indik, Lega,
... see the book of Pismen (2006)]

• Spatial dynamics : the Swift-Hohenberg equation

[H. & Scheel, 2012]

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THE SWIFT-HOHENBERG EQUATION

• Swift-Hohenberg equation

$$\mathsf{u}_{\mathsf{t}} = -(\mathbf{\Delta}+1)^2\mathsf{u} + \mu\mathsf{u} - \mathsf{u}^3$$

• grain boundaries : steady solutions



• anisotropic Swift-Hohenberg equation

$$u_{t} = -(\Delta + 1)^{2}u + \mu u - u^{3} + \beta u_{xx}$$

• dislocations : traveling waves



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SPATIAL DYNAMICS

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1 Dynamical system

$$\frac{\mathsf{dU}}{\mathsf{dx}} = \mathcal{A}(\mu,\mathsf{k},\mathsf{c},\boldsymbol{\beta})\mathsf{U} + \mathcal{F}(\mathsf{U})$$

- rolls \leftrightarrow equilibria
- dislocations / grain boundaries \leftrightarrow heteroclinic orbits

SPATIAL DYNAMICS

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1 Dynamical system

$$\frac{\mathsf{dU}}{\mathsf{dx}} = \mathcal{A}(\mu, \mathsf{k}, \mathsf{c}, \beta)\mathsf{U} + \mathcal{F}(\mathsf{U})$$

- rolls \leftrightarrow equilibria
- dislocations / grain boundaries \leftrightarrow heteroclinic orbits
- **2** Parameters : equation : μ , β
 - *y-periodic solutions :* wavenumber *k*
 - traveling waves : speed c
 - bifurcation points : co-existence of rolls with dislocations : different wavenumbers grain boundaries : different orientations
 - dispersion relation

REDUCED DYNAMICS

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1 Dynamical system

$$\frac{\mathsf{d}\mathsf{U}}{\mathsf{d}\mathsf{x}} = \mathcal{A}(\mu,\mathsf{k},\mathsf{c},\boldsymbol{\beta})\mathsf{U} + \mathcal{F}(\mathsf{U})$$

2 Parameters :

3 Center manifold reduction

4 Reduced system : find a heteroclinic orbit

REDUCED DYNAMICS

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2 Parameters :

3 Center manifold reduction

④ Reduced system : find a heteroclinic orbit dislocations : ODE in ℝ⁴
grain boundaries : ODE in ℝ¹² !

EXISTENCE OF GRAIN BOUNDARIES

• Reduced system : ODE in \mathbb{R}^{12}

$$\begin{array}{rcl} A'_{0} & = & \mathrm{i}A_{0} + B_{0} - \frac{\mathrm{i}}{4} \left(\mu a_{0} - a_{0} (a_{0}^{2} + 6a_{+}\overline{a_{+}}) \right) + \dots \\ B'_{0} & = & \mathrm{i}B_{0} - \frac{1}{4} \left(\mu a_{0} - a_{0} (a_{0}^{2} + 6a_{+}\overline{a_{+}}) \right) + \dots \\ A'_{+} & = & \mathrm{i}k_{x}A_{+} + B_{+} - \frac{\mathrm{i}}{4k_{x}^{3}} \left(\mu a_{+} - 3a_{+} (a_{0}^{2} + a_{+}\overline{a_{+}}) \right) + \dots \\ B'_{+} & = & \mathrm{i}k_{x}B_{+} - \frac{1}{4k_{x}^{2}} \left(\mu a_{+} - 3a_{+} (a_{0}^{2} + a_{+}\overline{a_{+}}) \right) + \dots \\ A'_{-} & = & \mathrm{i}k_{x}A_{-} + B_{-} - \frac{\mathrm{i}}{4k_{x}^{3}} \left(\mu \overline{a_{+}} - 3\overline{a_{+}} (a_{0}^{2} + a_{+}\overline{a_{+}}) \right) + \dots \\ B'_{-} & = & \mathrm{i}k_{x}B_{-} - \frac{1}{4k_{x}^{2}} \left(\mu \overline{a_{+}} - 3\overline{a_{+}} (a_{0}^{2} + a_{+}\overline{a_{+}}) \right) + \dots \end{array}$$

 $a_0=A_0+\overline{A_0},\ b_0=B_0+\overline{B_0},\ a_+=A_++\overline{A_-},\ a_-=A_+-\overline{A_-},\ b_+=B_++\overline{B_-}$

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Normal form and scalings

LEADING ORDER SYSTEM

$$\begin{array}{rcl} C_0'' & = & -\frac{1}{4}C_0 + \frac{3}{4}C_0(|C_0|^2 + 2|C_+|^2 + 2|C_-|^2) \\ C_+'' & = & -\frac{1}{4k_x^2}C_+ + \frac{3}{4k_x^2}C_+(2|C_0|^2 + |C_+|^2 + 2|C_-|^2) \\ C_-'' & = & -\frac{1}{4k_x^2}C_- + \frac{3}{4k_x^2}C_-(2|C_0|^2 + 2|C_+|^2 + |C_-|^2) \end{array}$$

- existence of a heteroclinic orbit $(0, C^{\star}_{+}, C^{\star}_{-})$
 - $(C_0^\star,C_+^\star,0)$?

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[van den Berg & van der Vorst, 1997]

- persistence of the heteroclinic orbit
 - analysis of the linearized operator
 - *implicit function theorem*

existence of grain boundaries

More Defects

• Some may be treatable by related methods ...









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More Defects

• Some may be treatable by related methods ...









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• Some cannot be treated by any of these methods ...





THIRD EXAMPLE

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LE PROBLÈME PHYSIQUE



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Résonateur optique



- Permet de générer des horloges de ultra-haute précision.
- Applications en aérospatiale, en télécommunications (radars, GPS...).
LE PROBLÈME MATHÉMATIQUE

(Lm^B)

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Équation de Lugiato-Lefever

$$\frac{\partial \psi}{\partial t} = -(1 + i\alpha)\psi + i\psi |\psi|^2 - i\beta \frac{\partial^2 \psi}{\partial x^2} + F$$

 But : étudier la dynamique des solutions en fonction des paramètres α ∈ ℝ et F > 0.

L'ÉQUATION DE LUGIATO-LEFEVER

$$\frac{\partial \psi}{\partial t} = -\left(1 + i\alpha\right)\psi + i\psi\left|\psi\right|^{2} - i\beta\frac{\partial^{2}\psi}{\partial x^{2}} + F$$

• Solutions les plus simples : les équilibres.

L'ÉQUATION DE LUGIATO-LEFEVER

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Stabilité des équilibres

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• Stabilité en tant que solutions de l'équation différentielle

$$\frac{\partial \psi}{\partial t} = -\left(1 + i\alpha\right)\psi + i\psi\left|\psi\right|^{2} + F$$

• On étudie les valeurs propres de la matrice de l'équation linéarisée.

Stabilité des équilibres

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ÉTUDE DES BIFURCATIONS

- **Bifurcations :** changement qualitatif dans la dynamique des solutions lorsque les paramètres varient.
 - On se restreint aux solutions stationnaires.
 - On étudie les valeurs propres de la matrice de l'équation linéarisée.

ÉTUDE DES BIFURCATIONS

- **Bifurcations :** changement qualitatif dans la dynamique des solutions lorsque les paramètres varient.
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LA SUITE ...



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• Calcul des formes normales : existence des solutions.



LA SUITE ...



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• Calcul des formes normales : existence des solutions.



- Étude de la stabilité des solutions périodiques.
- Interprétation physique des résultats mathématiques.

NONLINEAR WAVES AND PATTERNS





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